

Oscillations

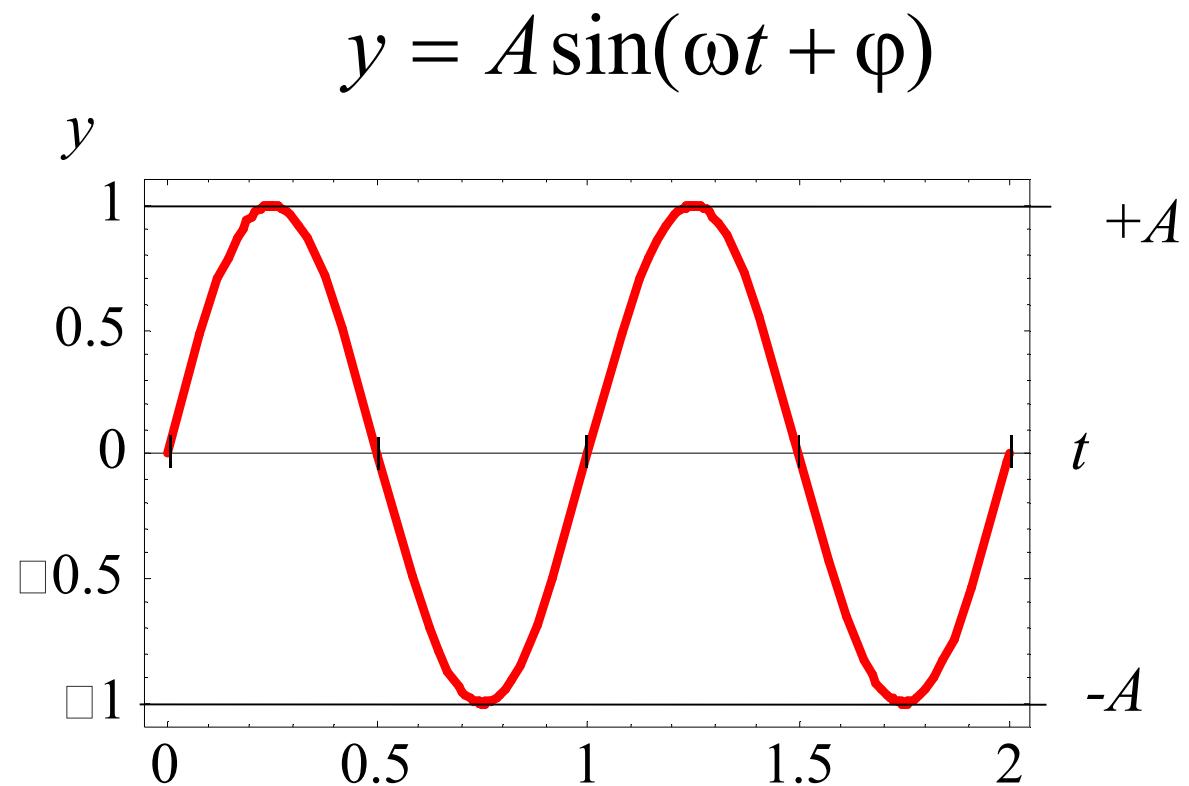
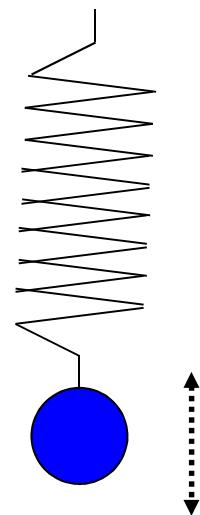
Free non-damped oscillations.

Superposition of oscillations, beats.

Free damped oscillations. Resonance.

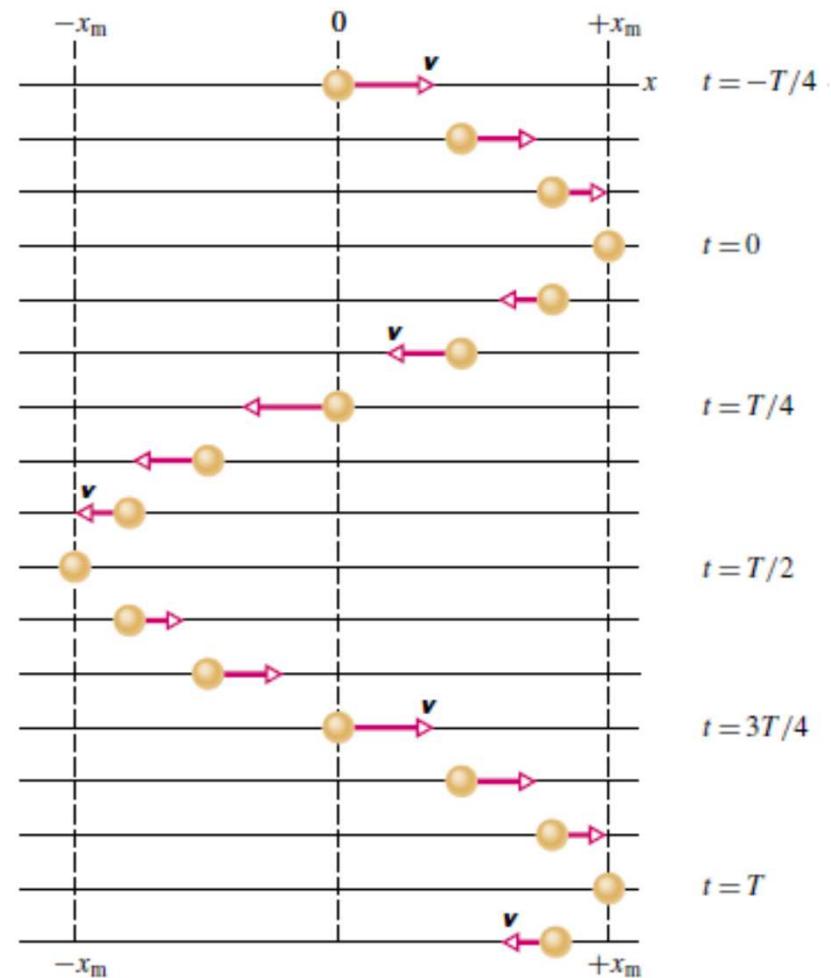
Harmonic motion

Periodic and space limited motion



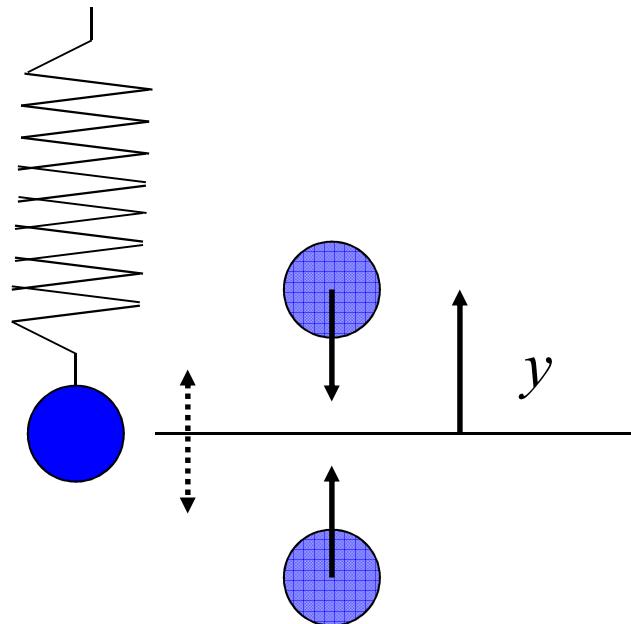
Harmonic motion

Displacement and
velocity change
magnitude and
direction



Restoring force

(Linear) restoring force



$$F = -ky$$

k ... spring constant
 y ... displacement

Type of oscillations

- Periodic
 - Non-periodic
 - Linear
 - Non-linear
- Free
 - etc.
- Forced
- Damped
- Non-damped

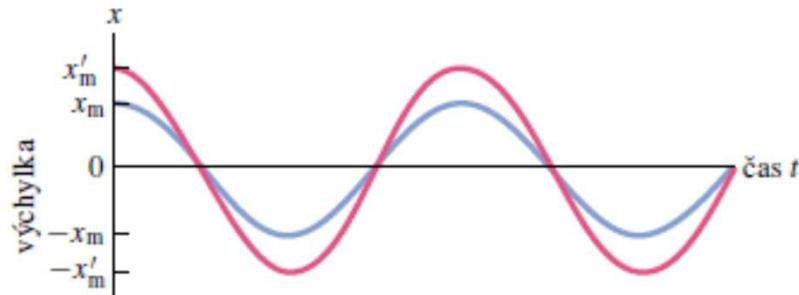
Non-damped oscillation motion parameters

Equation of motion solution $y = A \sin(\omega t + \varphi)$

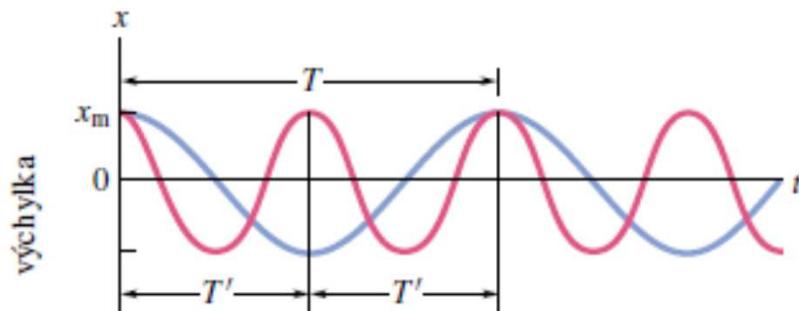
- Displacement - y
- Amplitude - A
- Angular frequency – ω , frequency - f
- Period – $\omega = \frac{2\pi}{T} = 2\pi f$
- Initial phase - φ

Oscillation parameters

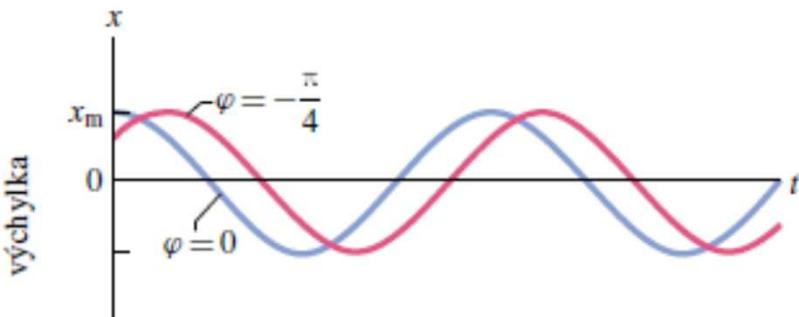
Amplitude



Period



Initial phase



Equation of motion for harmonic oscillations

Equation of motion for restoring force

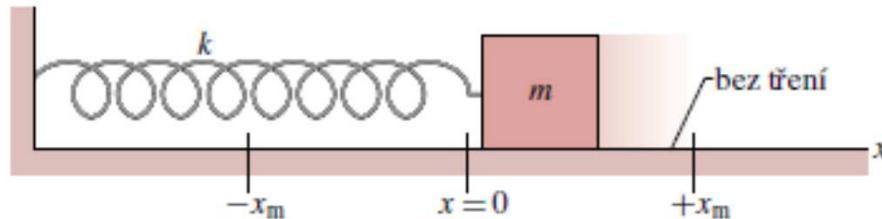
$$m \frac{d^2 y}{dt^2} = -ky$$

$$\frac{d^2 y}{dt^2} + \omega_0^2 y = 0$$

2nd order ordinary differential equation

Angular frequency $\omega_0^2 = \frac{k}{m}$

Angular frequency and oscillation period



- Angular frequency

$$\omega = \sqrt{\frac{k}{m}}$$

- Period

$$T = 2\pi\sqrt{\frac{m}{k}}$$

General solution of equation of motion

General solution $y = A \sin(\omega t + \varphi)$

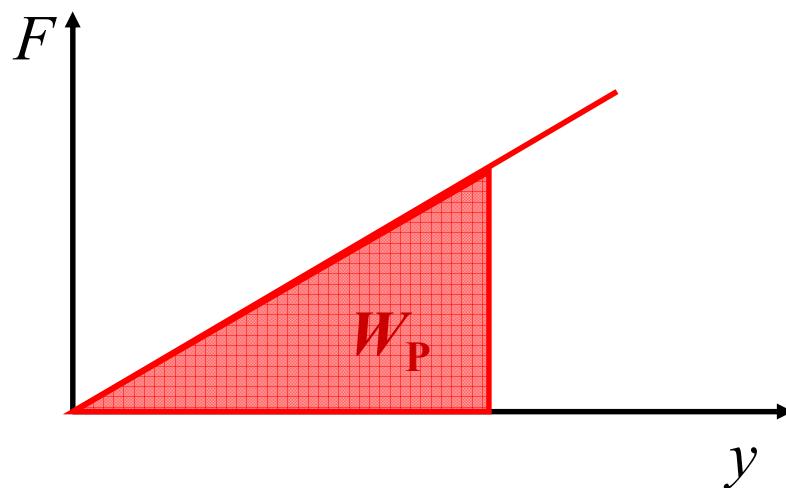
Two parameters – A, φ

Two initial conditions necessary, e.g.
maximum displacement and zero velocity

$$\begin{aligned} y(0) &= A \\ y'(0) &= 0 \end{aligned} \quad \longrightarrow \quad y = A \sin\left(\omega t + \frac{\pi}{2}\right) = A \cos(\omega t)$$

Potential energy of stressed spring

Restoring harmonic force $F = -ky$



$$W_P = - \int_0^{y_0} F dy = \int_0^{y_0} ky dy = \frac{1}{2} k y_0^2$$

Derivation of equation of motion from energy balance

- Kinetic energy of mass

$$W_K = \frac{1}{2}mv^2$$

- Potential energy of deformed spring

$$W_P = \frac{1}{2}ky^2$$

Energy conservation law – isolated system
mass + spring

$$\frac{1}{2}mv^2 + \frac{1}{2}ky^2 = \text{konst.}$$

Velocity and acceleration for oscillatory motion

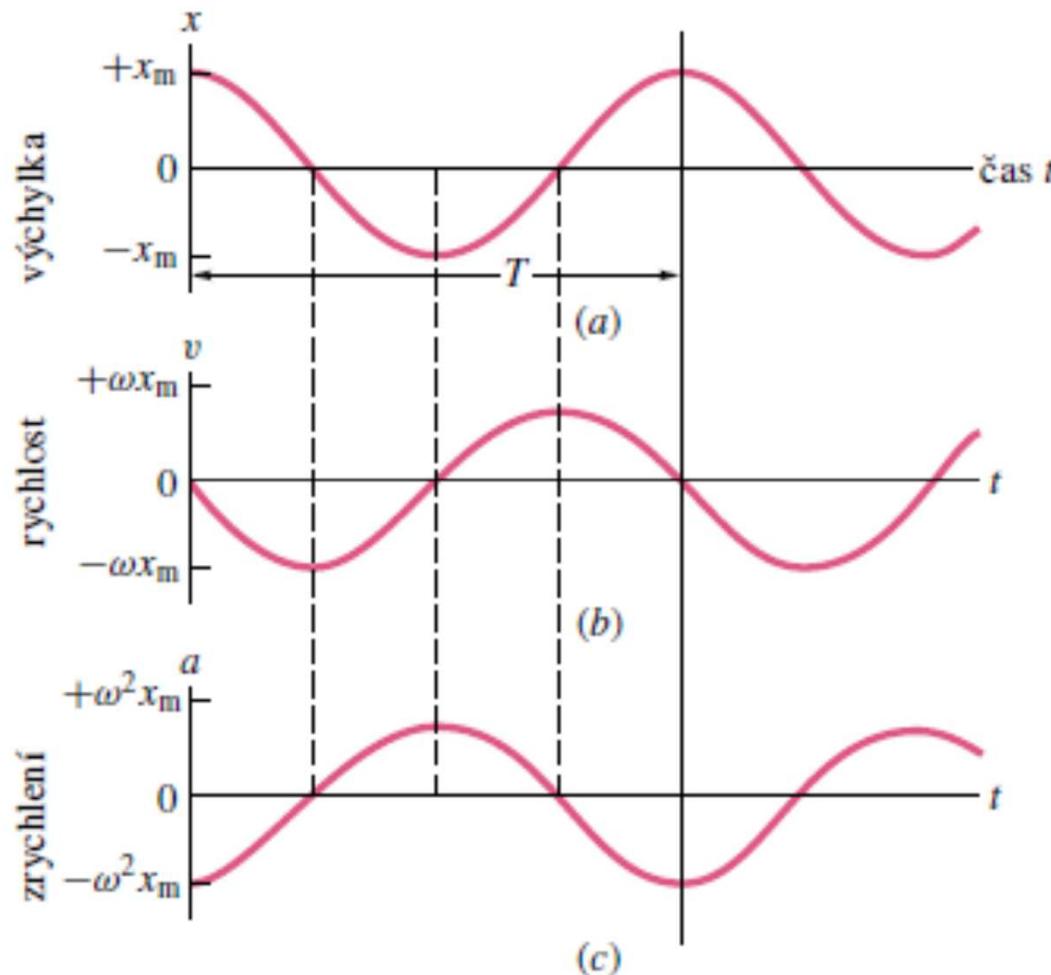
- Displacement $y = A \sin(\omega t + \varphi)$
- Velocity

$$v = \frac{dy}{dt} = A\omega \cos(\omega t + \varphi) = A\omega \sin(\omega t + \varphi + \frac{\pi}{2})$$

- Acceleration $a = -\omega^2 y$

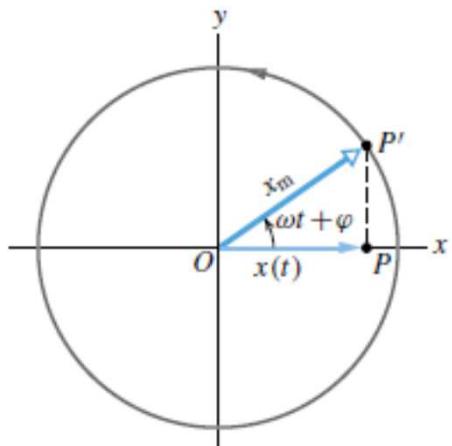
$$a = \frac{dv}{dt} = -A\omega^2 \sin(\omega t + \varphi) = A\omega^2 \sin(\omega t + \varphi + \pi)$$

Displacement, velocity and acceleration

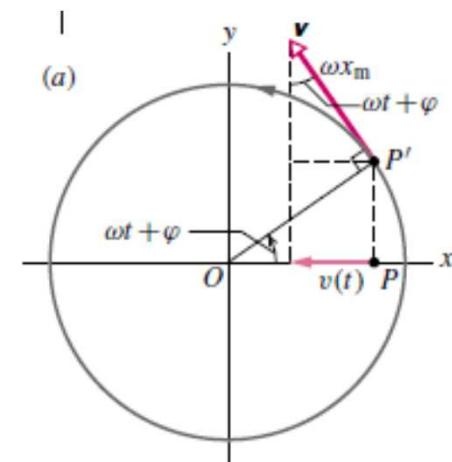


Circular motion analogy

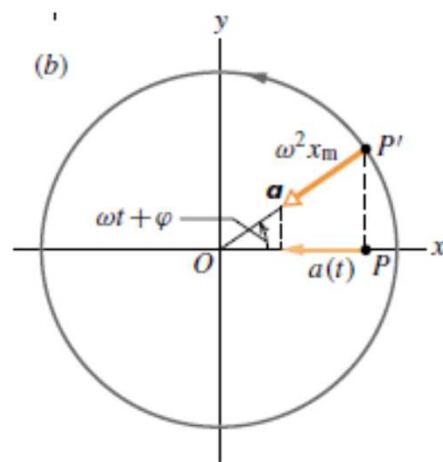
Displacement



Velocity



Acceleration



Energy of oscillation motion

Doubled frequency for energy changes!

- Kinetic

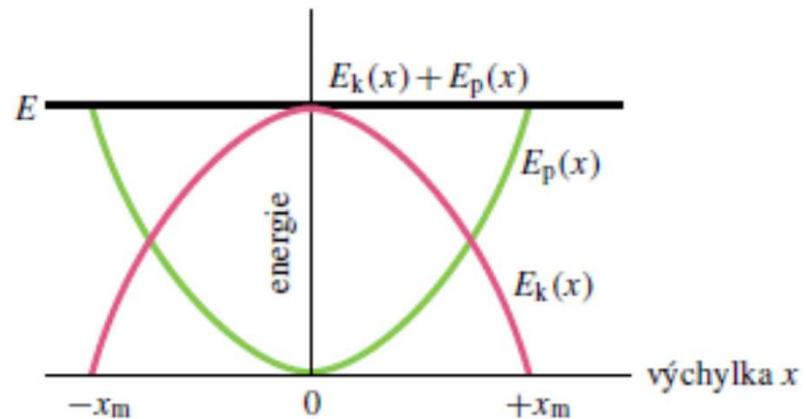
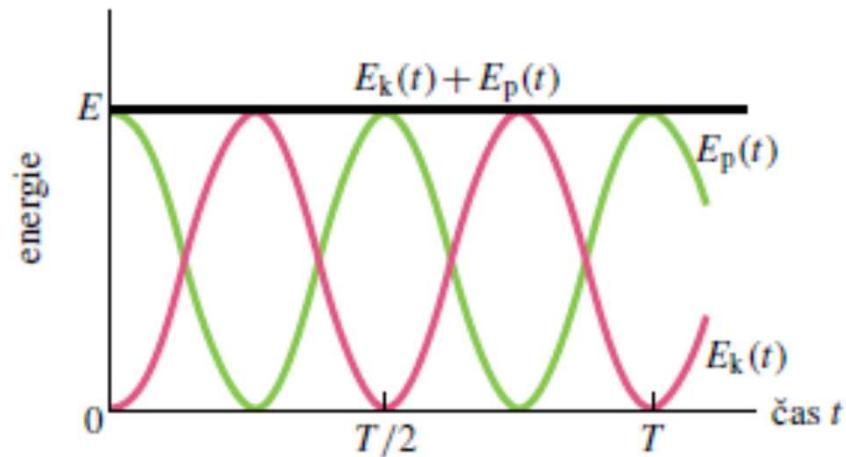
$$W_K = \frac{1}{2}mv^2 = \frac{1}{2}m\omega^2 A^2 \sin^2(\omega t) = \frac{1}{4}m\omega^2 A^2(1 - \cos(2\omega t))$$

- Potential

$$W_P = \frac{1}{2}ky^2 = \frac{1}{2}kA^2 \cos^2(\omega t) = \frac{1}{4}kA^2(1 + \cos(2\omega t))$$

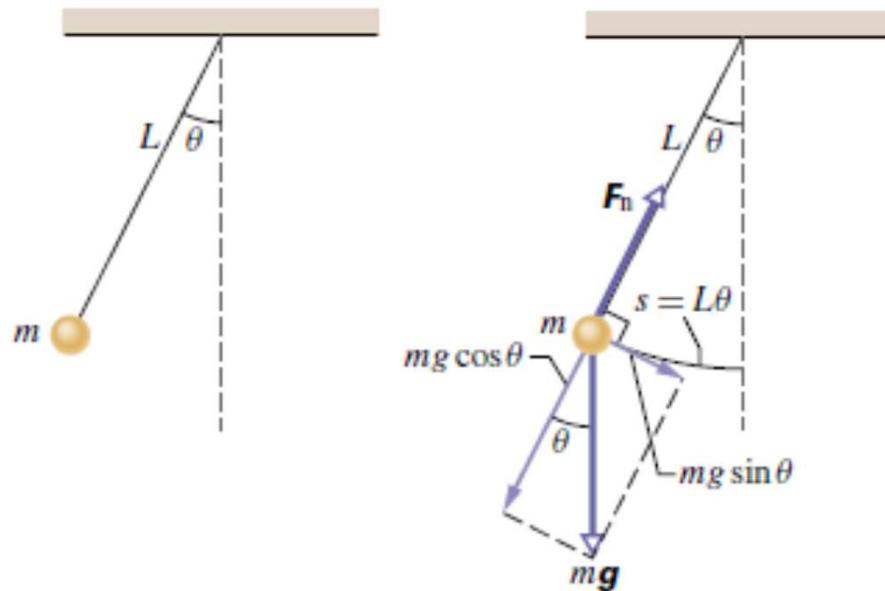
Displacement $y = A \cos(\omega t)$

Energy of oscillatory motion



Simple pendulum

Point mass (pendulum bob) and lightweight cord

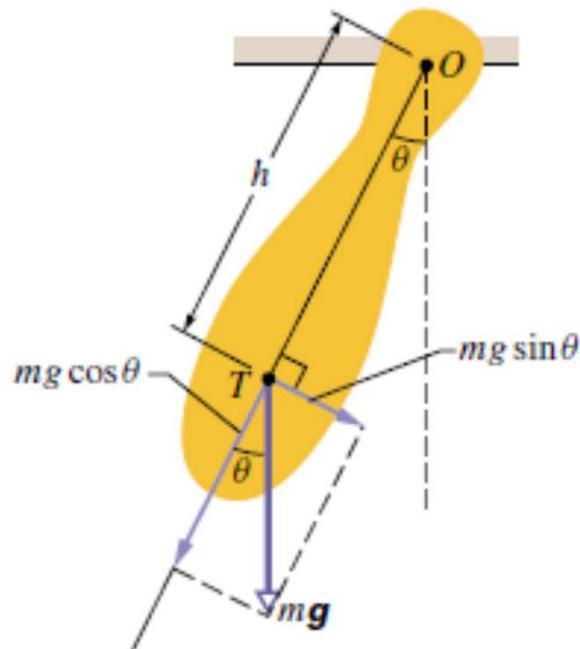


Period

$$T = 2\pi \sqrt{\frac{L}{g}}$$

Physical pendulum

Rigid body rotating by the axis (off center of mass)



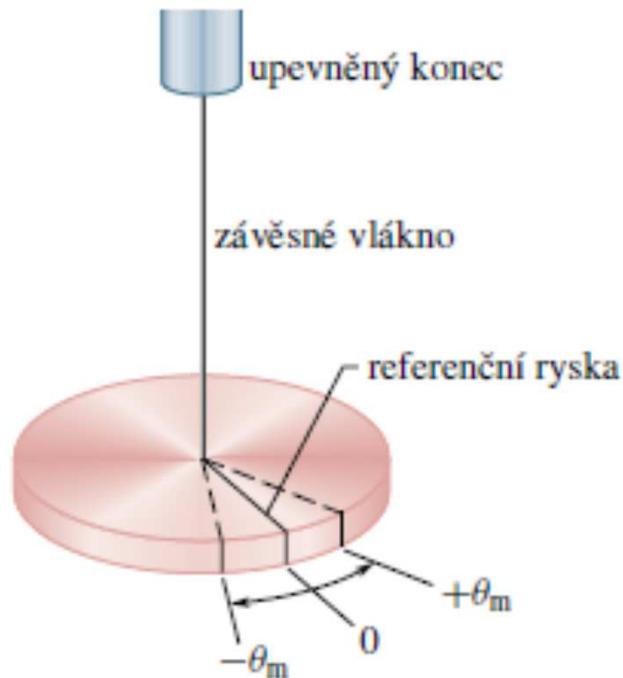
Period

$$T = 2\pi \sqrt{\frac{I}{mgh}}$$

inertial moment I

Torsion pendulum

Rigid body on lightweight cord deformed by torsion



Period

$$T = 2\pi \sqrt{\frac{I}{\kappa}}$$

inertial moment I
torsional toughness κ
torque (moment of force)

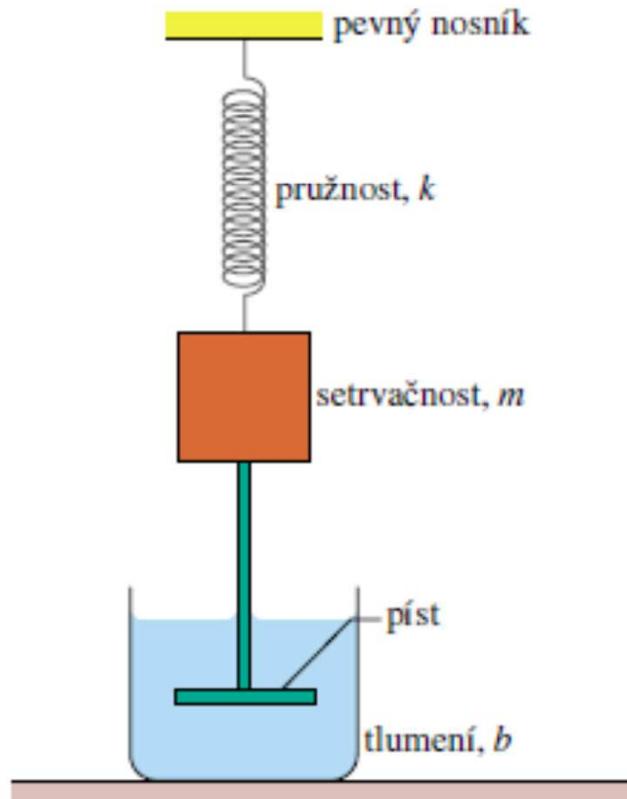
$$M = -\kappa\theta$$

Damped oscillations

Damping force

$$F_v = -kx - bv$$

Equation of motion



$$m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

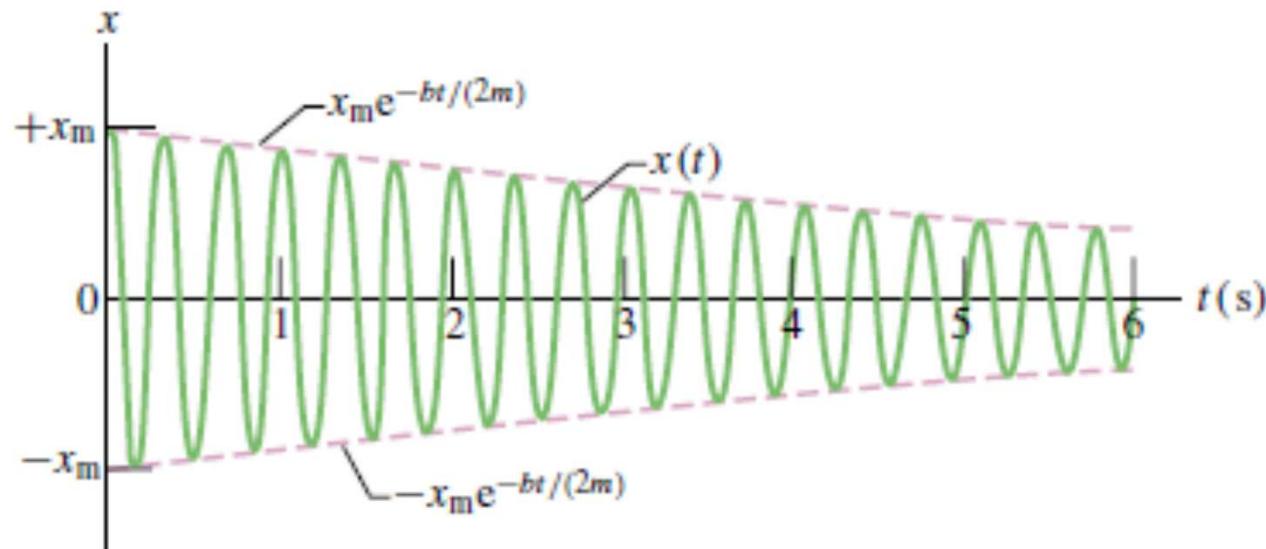
$$x(t) = x_m e^{-bt/(2m)} \cos(\omega' t + \varphi)$$

$$\omega' = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$$

Damped oscillations

exponential amplitude decay

$$x(t) = x_m e^{-bt/(2m)} \cos(\omega' t + \varphi)$$
$$\omega' = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}.$$



Attenuation and logarithmic decrement

Amplitude decays after every period by

$$\text{Attenuation } \frac{\frac{x_m e^{-\frac{bt}{2m}}}{x_m e^{-\frac{b(t+T)}{2m}}} = e^{\frac{bT}{2m}}}$$

Logarithmic decrement [Bell, dB]

$$\log \frac{x(t)}{x(t+T)} = \frac{bT}{2m}$$

Example

Logarithmic decrement 1dB (0.1B) means decay
of oscillation amplitude per cycle

$$\frac{x(t)}{x(t+T)} = 10^{0.1} = 1.259$$

3dB corresponds to the amplitude decay per
cycle approx. to one half

$$\frac{x(t)}{x(t+T)} = 10^{0.3} = 1.995$$

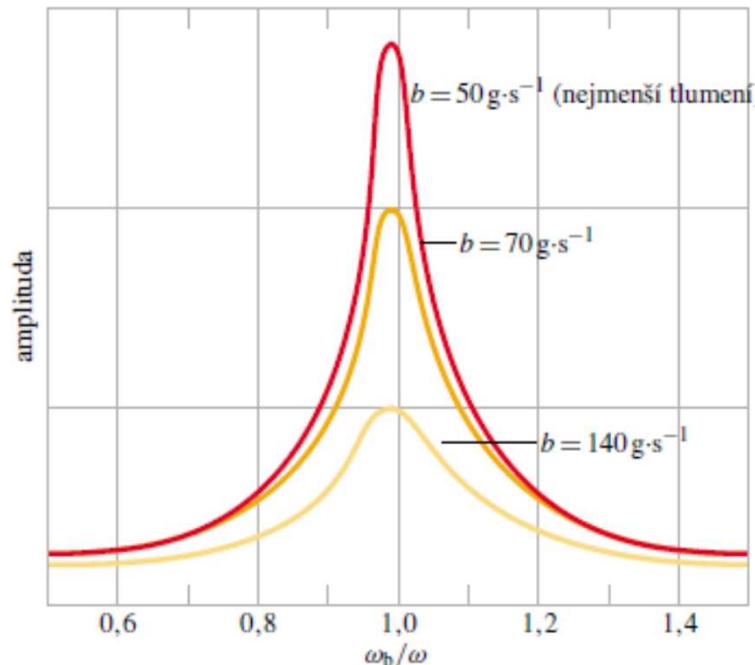
Forced vibrations

Oscillation is forced by periodic force at the frequency ω_b

$$x(t) = x_m \cos(\omega_b t + \varphi)$$

Oscillation amplitude depends on frequency

Damping constant



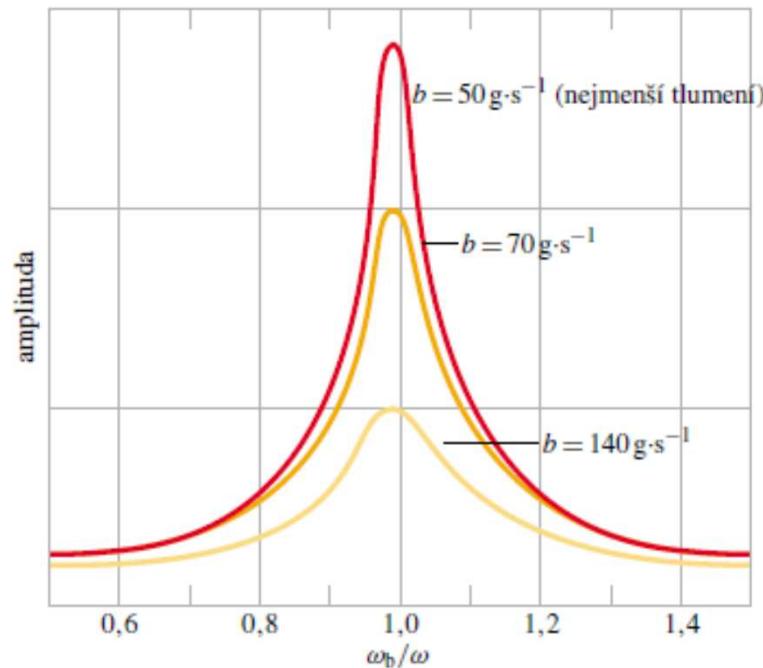
Resonance

Increase of amplitude
in the vicinity ω_b

$$x(t) = x_m \cos(\omega_b t + \varphi)$$

Resonance frequency

$$\omega_r = \sqrt{\omega_b^2 - \left(\frac{b}{2m}\right)^2}$$



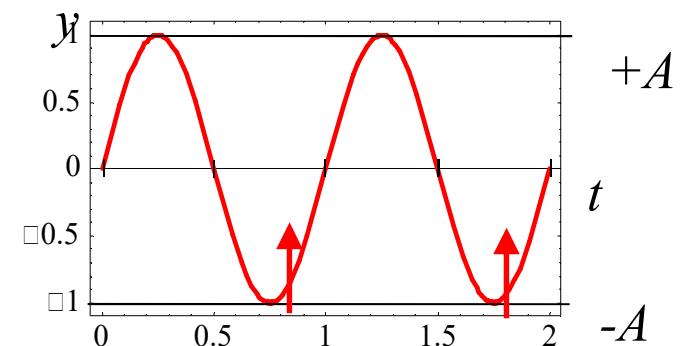
Parametric resonance

Different from resonance, change of oscillation parameters at the specific phase of oscillation

Example:

Swing – change in mass distribution in boundary displacement, keeps displacement amplitude without damping

Joo-joo (Maxwell's pendulum)



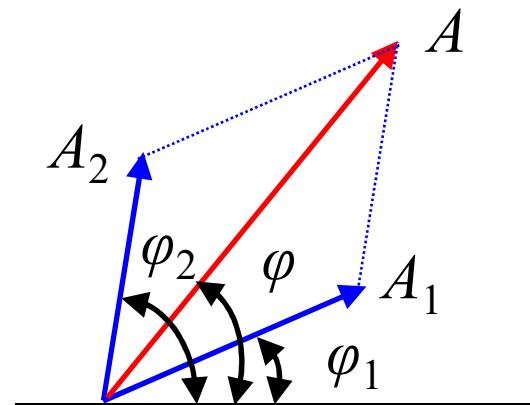
Superposition of oscillations with the same frequency

- Calculation

$$A^2 = A_1^2 + A_2^2 + 2A_1A_2 \cos(\varphi_2 - \varphi_1)$$

$$\tan \varphi = \frac{A_1 \sin \varphi_1 + A_2 \sin \varphi_2}{A_1 \cos \varphi_1 + A_2 \cos \varphi_2}$$

- Graphical method
by phase vectors



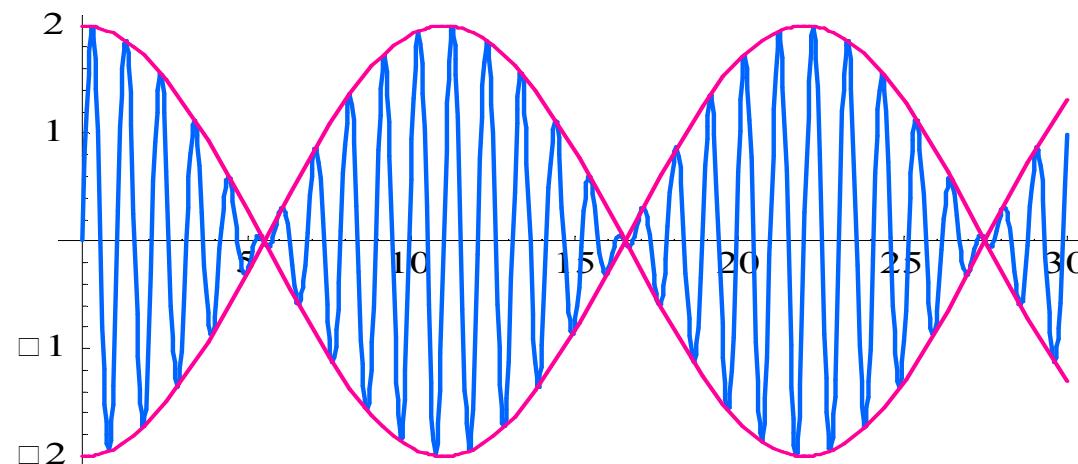
Beats

Oscillations of the same amplitude - superposition

$$y_1 = A \sin(\omega_1 t + \varphi_1) \quad y_2 = A \sin(\omega_2 t + \varphi_2)$$

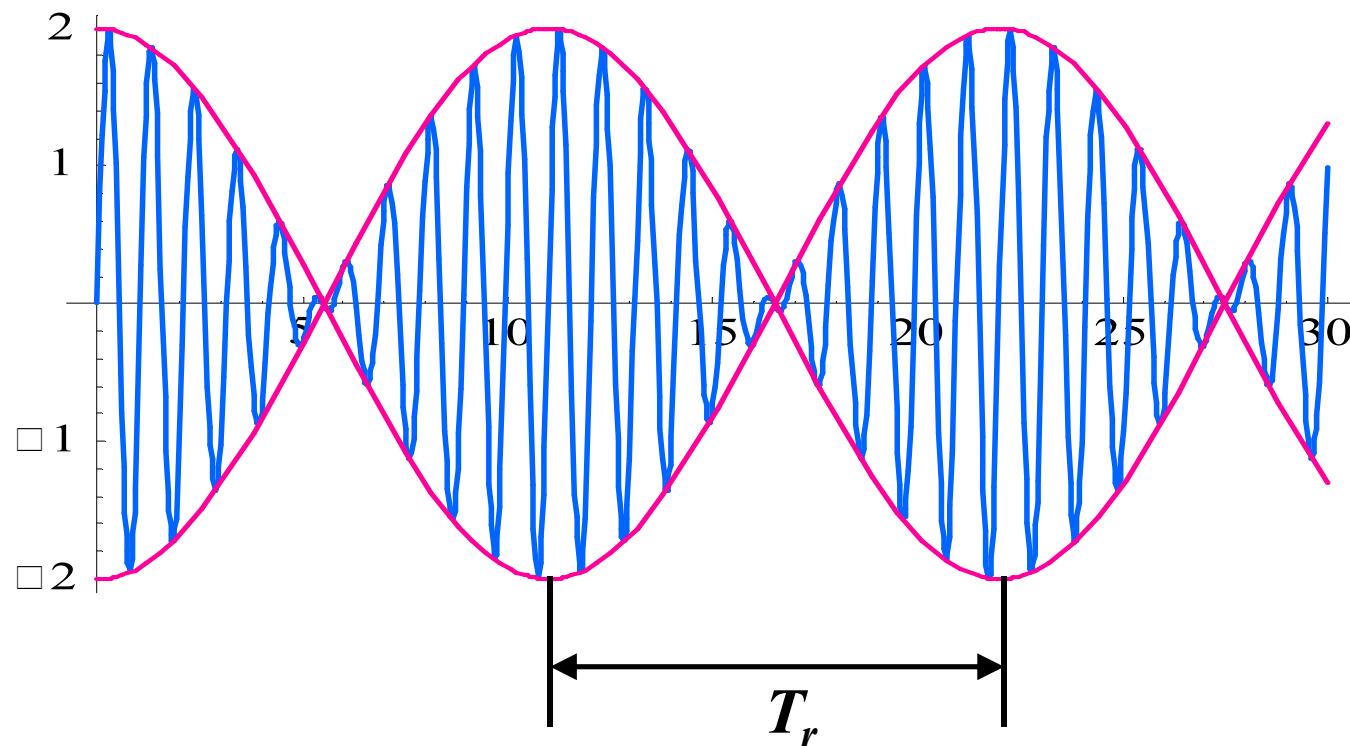
$$y = y_1 + y_2 =$$

$$= 2A \cos\left[\frac{1}{2}(\omega_1 - \omega_2)t + \frac{1}{2}(\varphi_1 - \varphi_2)\right] \sin\left[\frac{1}{2}(\omega_1 + \omega_2)t + \frac{1}{2}(\varphi_1 + \varphi_2)\right]$$



Beats

Period of beats $f_r = |f_1 - f_2|$, $T_r = 1/f_r$



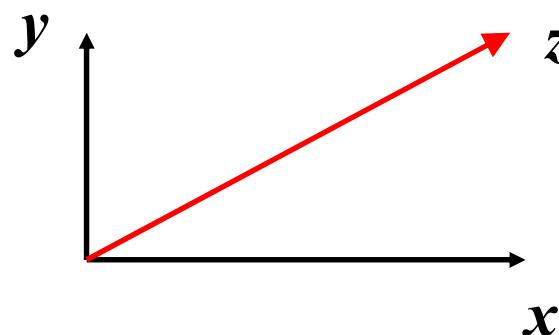
Lissajous pattern

Superposition of oscillations with different direction

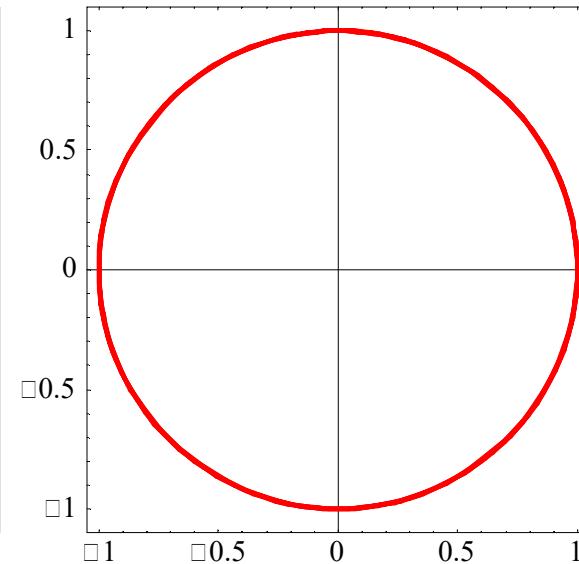
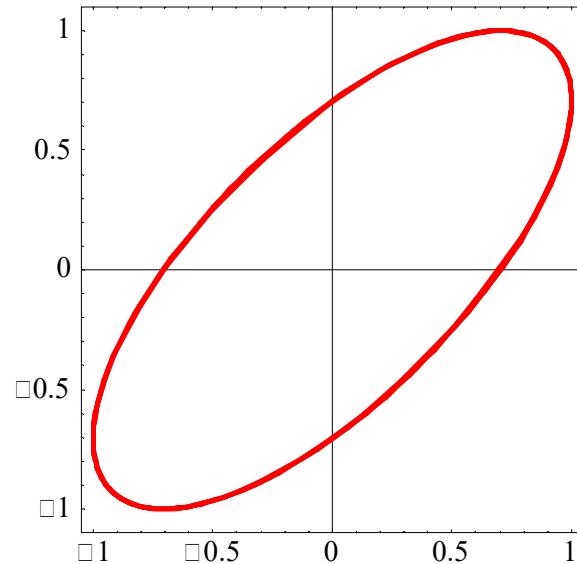
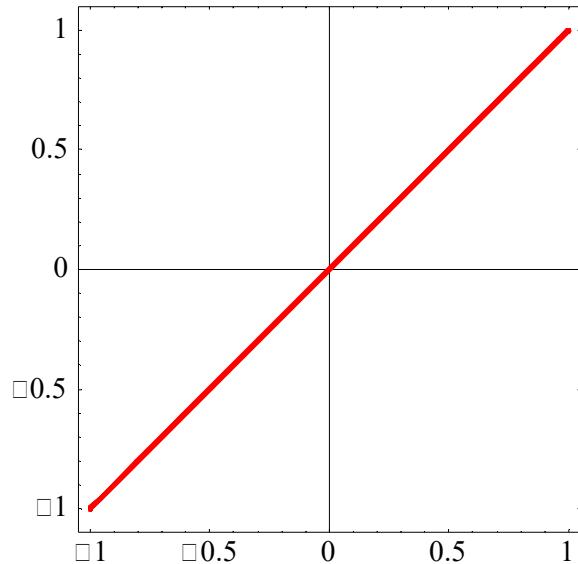
$$x = A_1 \sin(\omega_1 t + \varphi_1), y = A_2 \sin(\omega_2 t + \varphi_2)$$

$$\omega = 2\pi f$$

$$z = (x, y)$$



Lissajous pattern

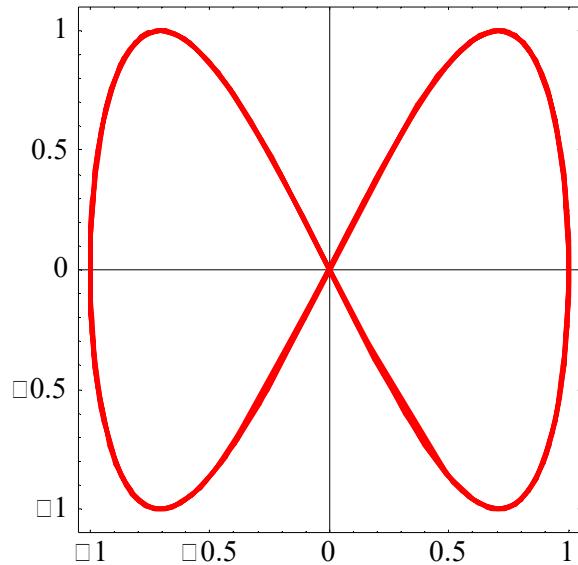


$f_1:f_2=1:1$
 $\varphi_1=0^\circ, \varphi_2=0^\circ$

$f_1:f_2=1:1$
 $\varphi_1=0^\circ, \varphi_2=45^\circ$

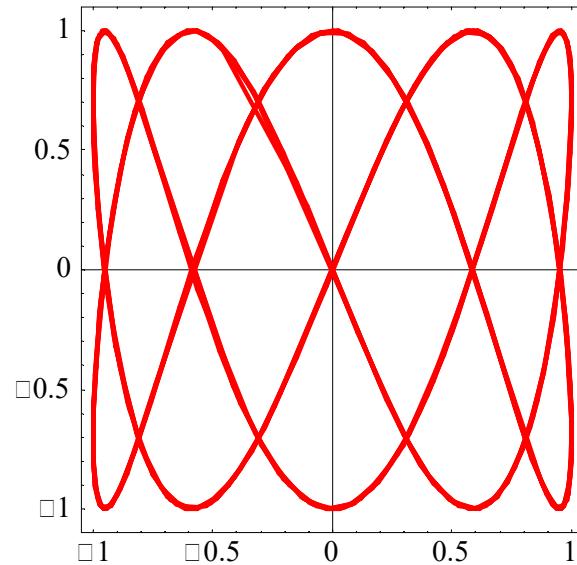
$f_1:f_2=1:1$
 $\varphi_1=0^\circ, \varphi_2=90^\circ$

Lissajous pattern



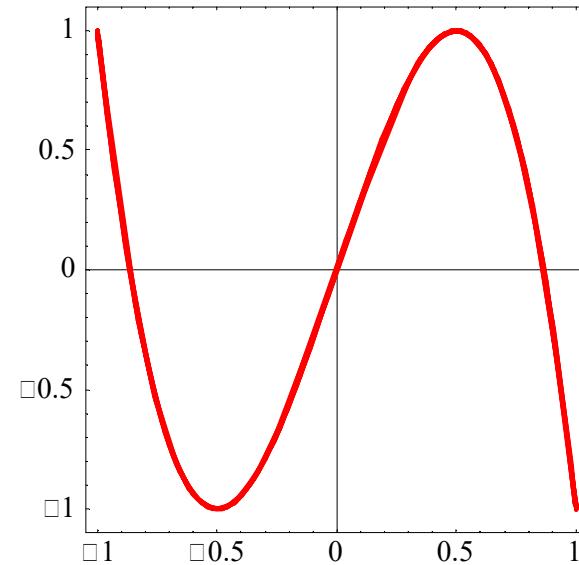
$$f_1:f_2=1:2$$

$$\varphi_1=0^\circ, \varphi_2=0^\circ$$



$$f_1:f_2=2:5$$

$$\varphi_1=0^\circ, \varphi_2=0^\circ$$



$$f_1:f_2=1:3$$

$$\varphi_1=0^\circ, \varphi_2=0^\circ$$

Literature

Pictures used from the book:

HALLIDAY, D., R. RESNICK, J. WALKER
Fyzika. Brno: VUTIUM, 2000. díl 2
Mechanika - Termodynamika